

Modelling a random rubber band

M. Hairer

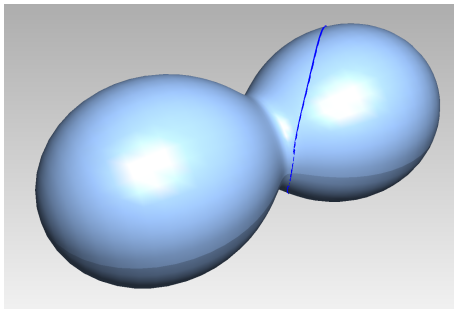
University of Warwick

Sendai, 03.09.2015

Joint work with L. Zambotti, Y. Bruned (Paris VI)

Introduction

Rubber band constrained to lie on a Manifold: $u: S^1 \rightarrow \mathcal{M}$.



Deterministic part of evolution: “length shortening” / heat flow.
Suitable L^2 -gradient flow for $\int g_u(\dot{u}, \dot{u}) dt$.

Local coordinates

In **local coordinates**, heat flow given by:

$$\partial_t u^\alpha = \partial_x^2 u^\alpha + \Gamma_{\beta\gamma}^\alpha(u) \partial_x u^\beta \partial_x u^\gamma =: (\Delta u)^\alpha .$$

Nonlinearity given by **Christoffel symbols** of Levi-Civita connection Γ .

Introduced by Eells-Sampson to show that every smooth map $\mathcal{N} \rightarrow \mathcal{M}$ is homotopic to a harmonic map.

Makes sense for any torsion-free connection (not just Levi-Civita): equilibria satisfy $\nabla_{\partial_x u} \partial_x u = 0$. Unit speed geodesics.

Local coordinates

In **local coordinates**, heat flow given by:

$$\partial_t u^\alpha = \partial_x^2 u^\alpha + \Gamma_{\beta\gamma}^\alpha(u) \partial_x u^\beta \partial_x u^\gamma =: (\Delta u)^\alpha .$$

Nonlinearity given by **Christoffel symbols** of Levi-Civita connection Γ .

Introduced by Eells-Sampson to show that every smooth map $\mathcal{N} \rightarrow \mathcal{M}$ is homotopic to a harmonic map.

Makes sense for any torsion-free connection (not just Levi-Civita): equilibria satisfy $\nabla_{\partial_x u} \partial_x u = 0$. Unit speed geodesics.

Local coordinates

In **local coordinates**, heat flow given by:

$$\partial_t u^\alpha = \partial_x^2 u^\alpha + \Gamma_{\beta\gamma}^\alpha(u) \partial_x u^\beta \partial_x u^\gamma =: (\Delta u)^\alpha .$$

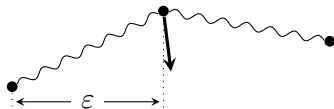
Nonlinearity given by **Christoffel symbols** of Levi-Civita connection Γ .

Introduced by Eells-Sampson to show that every smooth map $\mathcal{N} \rightarrow \mathcal{M}$ is homotopic to a harmonic map.

Makes sense for any torsion-free connection (not just Levi-Civita): equilibria satisfy $\nabla_{\partial_x u} \partial_x u = 0$. Unit speed geodesics.

Adding noise

Consider a discrete version:



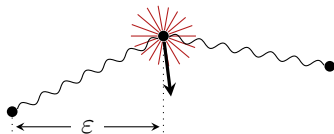
Add **independent noises** to each bead.

Basic question: how strong should the noise be to create an effect of order 1? Recall that noise of strength σ yields displacement $\mathcal{O}(\sigma\sqrt{t})$.

First guess: relevant timescale $\mathcal{O}(\epsilon^2)$, want displacement $\mathcal{O}(\epsilon) \Rightarrow \sigma \approx 1$.

Adding noise

Consider a discrete version:



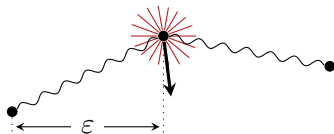
Add **independent noises** to each bead.

Basic question: how strong should the noise be to create an effect of order 1? Recall that noise of strength σ yields displacement $\mathcal{O}(\sigma\sqrt{t})$.

First guess: relevant timescale $\mathcal{O}(\varepsilon^2)$, want displacement $\mathcal{O}(\varepsilon) \Rightarrow \sigma \approx 1$.

Adding noise

Consider a discrete version:



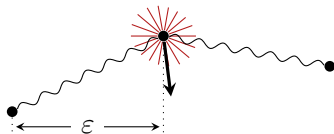
Add **independent noises** to each bead.

Basic question: how strong should the noise be to create an effect of order 1? Recall that noise of strength σ yields displacement $\mathcal{O}(\sigma\sqrt{t})$.

First guess: relevant timescale $\mathcal{O}(\varepsilon^2)$, want displacement $\mathcal{O}(\varepsilon) \Rightarrow \sigma \approx 1$.

Adding noise

Consider a discrete version:



Add **independent noises** to each bead.

Basic question: how strong should the noise be to create an effect of order 1? Recall that noise of strength σ yields displacement $\mathcal{O}(\sigma\sqrt{t})$.

First guess: relevant timescale $\mathcal{O}(\epsilon^2)$, want displacement $\mathcal{O}(\epsilon) \Rightarrow \sigma \approx 1$.

A convergence result

Continuous version:

$$\partial_t u_\varepsilon = \Delta u_\varepsilon + \sigma_i(u_\varepsilon) \xi_i^{(\varepsilon)},$$

ξ_i centred **Gaussian** with $\mathbf{E} \xi_i^{(\varepsilon)}(t, x) \xi_j^{(\varepsilon)}(s, y) = \varepsilon^{-2} \delta_{ij} \varrho(\frac{t-s}{\varepsilon^2}, \frac{y-x}{\varepsilon})$
for ϱ smooth compactly supported. (Factor ε^{-2} correct scaling for white noise in time.)

Theorem: As $\varepsilon \rightarrow 0$, $u_\varepsilon \rightarrow u$ where u solves

$$\partial_t u = \Delta u + c(\nabla_{\sigma_i} \sigma_i)(u),$$

for some c depending on ϱ .

Limit is **deterministic**, not really what we wanted...

A convergence result

Continuous version:

$$\partial_t u_\varepsilon = \Delta u_\varepsilon + \sigma_i(u_\varepsilon) \xi_i^{(\varepsilon)},$$

ξ_i centred **Gaussian** with $\mathbf{E} \xi_i^{(\varepsilon)}(t, x) \xi_j^{(\varepsilon)}(s, y) = \varepsilon^{-2} \delta_{ij} \varrho(\frac{t-s}{\varepsilon^2}, \frac{y-x}{\varepsilon})$
for ϱ smooth compactly supported. (Factor ε^{-2} correct scaling for white noise in time.)

Theorem: As $\varepsilon \rightarrow 0$, $u_\varepsilon \rightarrow u$ where u solves

$$\partial_t u = \Delta u + c(\nabla_{\sigma_i} \sigma_i)(u),$$

for some c depending on ϱ .

Limit is **deterministic**, not really what we wanted...

Stronger noise

To have any hope for **finite limit** with **stronger noise**, need $\nabla_{\sigma_i} \sigma_i = 0$, or need to compensate this term.

Geometric interpretation: If vectors σ_i span the whole tangent space at each point, then they define an (inverse) metric by

$$g^{\alpha\beta}(u) = \sigma_i^\alpha(u) \sigma_i^\beta(u) .$$

Metric + connection \Rightarrow Laplace-Beltrami operator Δ .

Fact: diffusion $\dot{x} = \sigma_i(x) \xi_i$ with i.i.d. white noises ξ_i interpreted in Stratonovich sense has generator Δ **if and only if** $\nabla_{\sigma_i} \sigma_i = 0$.

Yields natural (nonlinear) centering condition, should allow to go from LLN to CLT.

Stronger noise

To have any hope for **finite limit** with **stronger noise**, need $\nabla_{\sigma_i} \sigma_i = 0$, or need to compensate this term.

Geometric interpretation: If vectors σ_i span the whole tangent space at each point, then they define an (inverse) metric by

$$g^{\alpha\beta}(u) = \sigma_i^\alpha(u) \sigma_i^\beta(u) .$$

Metric + connection \Rightarrow Laplace-Beltrami operator Δ .

Fact: diffusion $\dot{x} = \sigma_i(x) \xi_i$ with i.i.d. white noises ξ_i interpreted in Stratonovich sense has generator Δ **if and only if** $\nabla_{\sigma_i} \sigma_i = 0$.

Yields natural (nonlinear) centering condition, should allow to go from LLN to CLT.

Stronger noise

To have any hope for **finite limit** with **stronger noise**, need $\nabla_{\sigma_i} \sigma_i = 0$, or need to compensate this term.

Geometric interpretation: If vectors σ_i span the whole tangent space at each point, then they define an (inverse) metric by

$$g^{\alpha\beta}(u) = \sigma_i^\alpha(u) \sigma_i^\beta(u) .$$

Metric + connection \Rightarrow Laplace-Beltrami operator Δ .

Fact: diffusion $\dot{x} = \sigma_i(x) \xi_i$ with i.i.d. white noises ξ_i interpreted in Stratonovich sense has generator Δ **if and only if** $\nabla_{\sigma_i} \sigma_i = 0$.

Yields natural (nonlinear) centering condition, should allow to go from LLN to CLT.

Stronger noise

To have any hope for **finite limit** with **stronger noise**, need $\nabla_{\sigma_i} \sigma_i = 0$, or need to compensate this term.

Geometric interpretation: If vectors σ_i span the whole tangent space at each point, then they define an (inverse) metric by

$$g^{\alpha\beta}(u) = \sigma_i^\alpha(u) \sigma_i^\beta(u) .$$

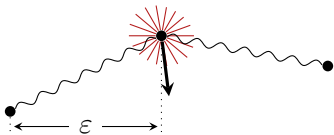
Metric + connection \Rightarrow Laplace-Beltrami operator Δ .

Fact: diffusion $\dot{x} = \sigma_i(x) \xi_i$ with i.i.d. white noises ξ_i interpreted in Stratonovich sense has generator Δ **if and only if** $\nabla_{\sigma_i} \sigma_i = 0$.

Yields natural (nonlinear) centering condition, should allow to go from LLN to CLT.

How strong?

Recall discrete picture:

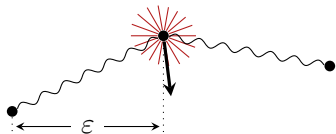


Random limit \Rightarrow expect solutions to behave like random walks
 \Rightarrow each “step” should be of size $\mathcal{O}(\sqrt{\varepsilon})$ to yield effect $\mathcal{O}(1)$ after $\mathcal{O}(1/\varepsilon)$ steps.

Conclusion: want $\sigma \approx \varepsilon^{-1/2}$ to have displacement $\mathcal{O}(\sqrt{\varepsilon})$ over timescale $\mathcal{O}(\varepsilon^2) \Rightarrow$ want $\mathbf{E}\xi_i^{(\varepsilon)}(t, x)\xi_i^{(\varepsilon)}(s, y) = \varepsilon^{-3}\varrho(\frac{t-s}{\varepsilon^2}, \frac{y-x}{\varepsilon})$.
Yields space-time white noise in the limit.

How strong?

Recall discrete picture:



Random limit \Rightarrow expect solutions to behave like random walks
 \Rightarrow each “step” should be of size $\mathcal{O}(\sqrt{\varepsilon})$ to yield effect $\mathcal{O}(1)$ after $\mathcal{O}(1/\varepsilon)$ steps.

Conclusion: want $\sigma \approx \varepsilon^{-1/2}$ to have displacement $\mathcal{O}(\sqrt{\varepsilon})$ over timescale $\mathcal{O}(\varepsilon^2) \Rightarrow$ want $\mathbf{E}\xi_i^{(\varepsilon)}(t, x)\xi_i^{(\varepsilon)}(s, y) = \varepsilon^{-3}\varrho\left(\frac{t-s}{\varepsilon^2}, \frac{y-x}{\varepsilon}\right)$.
Yields **space-time white noise** in the limit.

Main result

Consider similarly to before

$$\partial_t u_\varepsilon = \Delta u_\varepsilon + \sigma_i(u_\varepsilon) \xi_i^{(\varepsilon)} + h(u_\varepsilon) - \frac{\log \varepsilon}{4\pi\sqrt{3}} \hat{h}(u_\varepsilon),$$

but now $\mathbf{E} \xi_i^{(\varepsilon)}(t, x) \xi_i^{(\varepsilon)}(s, y) = \varepsilon^{-3} \varrho\left(\frac{t-s}{\varepsilon^2}, \frac{y-x}{\varepsilon}\right)$ and $\nabla_{\sigma_i} \sigma_i = 0$. For R curvature tensor built from Γ , set $\hat{h}^\alpha = R_{\beta\gamma\delta}^\alpha g^{\gamma\eta} \nabla_\eta g^{\delta\beta}$.

Theorem: There is a choice of vector field h depending on ϱ such that, as $\varepsilon \rightarrow 0$, $u_\varepsilon \rightarrow u$ for a limiting process u independent of ϱ . Comes with **intrinsic characterisation** of limiting process, morally

$$\partial_t u = \Delta u + \sigma_i(u) \xi_i.$$

Main result

Consider similarly to before

$$\partial_t u_\varepsilon = \Delta u_\varepsilon + \sigma_i(u_\varepsilon) \xi_i^{(\varepsilon)} + h(u_\varepsilon) - \frac{\log \varepsilon}{4\pi\sqrt{3}} \hat{h}(u_\varepsilon),$$

but now $\mathbf{E} \xi_i^{(\varepsilon)}(t, x) \xi_i^{(\varepsilon)}(s, y) = \varepsilon^{-3} \varrho\left(\frac{t-s}{\varepsilon^2}, \frac{y-x}{\varepsilon}\right)$ and $\nabla_{\sigma_i} \sigma_i = 0$. For R curvature tensor built from Γ , set $\hat{h}^\alpha = R_{\beta\gamma\delta}^\alpha g^{\gamma\eta} \nabla_\eta g^{\delta\beta}$.

Theorem: There is a choice of vector field h depending on ϱ such that, as $\varepsilon \rightarrow 0$, $u_\varepsilon \rightarrow u$ for a limiting process u **independent of ϱ** . Comes with **intrinsic characterisation** of limiting process, morally

$$\partial_t u = \Delta u + \sigma_i(u) \xi_i.$$

A few existing results

1. **Very large literature** on SPDEs in flat space. Even there, approximation is highly non-trivial if martingale structure is destroyed (cf. H. & Pardoux, '14).
2. Model of random string considered in (Funaki, '92), requires noise **smooth** in space.
3. Static case (i.e. make sense of the measure formally described by $\exp(-\int g_u(\dot{u}, \dot{u}) dt)$ “ Du ”) quite well understood, see Inoue & Maeda '85, Andersson & Driver '99. Different interpretations of “ Du ” yield slightly different measures.
4. Paraproduct-based theory developed by Gubinelli, Imkeller & Perkowski '14 does not seem to apply in this case.

A few existing results

1. **Very large literature** on SPDEs in flat space. Even there, approximation is highly non-trivial if martingale structure is destroyed (cf. H. & Pardoux, '14).
2. Model of random string considered in (Funaki, '92), requires noise **smooth** in space.
3. Static case (i.e. make sense of the measure formally described by $\exp(-\int g_u(\dot{u}, \dot{u}) dt)$ “ Du ”) quite well understood, see Inoue & Maeda '85, Andersson & Driver '99. Different interpretations of “ Du ” yield slightly different measures.
4. Paraproduct-based theory developed by Gubinelli, Imkeller & Perkowski '14 does not seem to apply in this case.

A few existing results

1. **Very large literature** on SPDEs in flat space. Even there, approximation is highly non-trivial if martingale structure is destroyed (cf. H. & Pardoux, '14).
2. Model of random string considered in (Funaki, '92), requires noise **smooth** in space.
3. Static case (i.e. make sense of the measure formally described by $\exp(-\int g_u(\dot{u}, \dot{u}) dt)$ “ Du ”) quite well understood, see Inoue & Maeda '85, Andersson & Driver '99. Different interpretations of “ Du ” yield slightly different measures.
4. Paraproduct-based theory developed by Gubinelli, Imkeller & Perkowski '14 does not seem to apply in this case.

A few existing results

1. **Very large literature** on SPDEs in flat space. Even there, approximation is highly non-trivial if martingale structure is destroyed (cf. H. & Pardoux, '14).
2. Model of random string considered in (Funaki, '92), requires noise **smooth** in space.
3. Static case (i.e. make sense of the measure formally described by $\exp(-\int g_u(\dot{u}, \dot{u}) dt)$ “ Du ”) quite well understood, see Inoue & Maeda '85, Andersson & Driver '99. Different interpretations of “ Du ” yield slightly different measures.
4. Paraproduct-based theory developed by Gubinelli, Imkeller & Perkowski '14 does not seem to apply in this case.

Main idea

One can “guess” the local behaviour of solutions:

$$\partial_t u^\alpha = \partial_x^2 u^\alpha + \Gamma_{\beta\gamma}^\alpha(u) \bullet \partial_x u^\beta \bullet \partial_x u^\gamma + \sigma_i^\alpha(u) \bullet \xi_i .$$

Near (x, t) , one would expect

$$u^\alpha(y, s) \approx u^\alpha(x, t) + \sigma_i^\alpha(u(x, t))(v_i(y, s) - v_i(x, t)) ,$$

with v_i solving $\partial_t v_i = \partial_x^2 v_i + \xi_i$.

Hope: Maybe defining the various products for v instead of u already yields enough information? Should be easier since v is explicit... **Unfortunately not:** Needs to go to higher order.

Main idea

One can “guess” the local behaviour of solutions:

$$\partial_t u^\alpha = \partial_x^2 u^\alpha + \Gamma_{\beta\gamma}^\alpha(u) \bullet \partial_x u^\beta \bullet \partial_x u^\gamma + \sigma_i^\alpha(u) \bullet \xi_i .$$

Near (x, t) , one would expect

$$u^\alpha(y, s) \approx u^\alpha(x, t) + \sigma_i^\alpha(u(x, t))(v_i(y, s) - v_i(x, t)) ,$$

with v_i solving $\partial_t v_i = \partial_x^2 v_i + \xi_i$.

Hope: Maybe defining the various products for v instead of u already yields enough information? Should be **easier** since v is explicit... **Unfortunately not:** Needs to go to higher order.

Main idea

One can “guess” the local behaviour of solutions:

$$\partial_t u^\alpha = \partial_x^2 u^\alpha + \Gamma_{\beta\gamma}^\alpha(u) \bullet \partial_x u^\beta \bullet \partial_x u^\gamma + \sigma_i^\alpha(u) \bullet \xi_i .$$

Near (x, t) , one would expect

$$u^\alpha(y, s) \approx u^\alpha(x, t) + \sigma_i^\alpha(u(x, t))(v_i(y, s) - v_i(x, t)) ,$$

with v_i solving $\partial_t v_i = \partial_x^2 v_i + \xi_i$.

Hope: Maybe defining the various products for v instead of u already yields enough information? Should be **easier** since v is explicit... **Unfortunately not:** Needs to go to higher order.

Higher order expansion

Graphical notations for symbols: \circ for noise, --- for heat kernel, --- for derivative of heat kernel. For example $(\Pi_{z_0}^{\circ})(z) = v(z) - v(z_0)$.
 (! Suitable recentering required !)

$$\begin{aligned}
 U = & u \mathbf{1} + \sigma \circ + \sigma \partial \sigma \circ \circ + \Gamma \sigma^2 \circ \circ + u' X \\
 & + 2\Gamma \sigma^2 \partial \sigma \circ \circ \circ + 2\Gamma^2 \sigma^3 \circ \circ \circ + \sigma^3 \partial \Gamma \circ \circ \circ \\
 & + \frac{1}{2} \sigma^2 \partial^2 \sigma \circ \circ \circ + \sigma (\partial \sigma)^2 \circ \circ \circ + \Gamma \sigma^2 \partial \sigma \circ \circ \circ \\
 & + u' \partial \sigma \circ + 2\Gamma \sigma u' \circ .
 \end{aligned}$$

Remark: Similar to expansions in Feynman diagrams. However, coefficients are not constant but depend on the solution itself.

Higher order expansion

Graphical notations for symbols: \circ for noise, --- for heat kernel, --- for derivative of heat kernel. For example $(\Pi_{z_0}^{\circ})(z) = v(z) - v(z_0)$.
 (! Suitable recentering required !) One has

$$\begin{aligned}
 U = & u \mathbf{1} + \sigma \circ + \sigma \partial \sigma \circ \circ + \Gamma \sigma^2 \circ \circ + u' X \\
 & + 2\Gamma \sigma^2 \partial \sigma \circ \circ + 2\Gamma^2 \sigma^3 \circ \circ + \sigma^3 \partial \Gamma \circ \circ \\
 & + \frac{1}{2} \sigma^2 \partial^2 \sigma \circ \circ + \sigma (\partial \sigma)^2 \circ \circ + \Gamma \sigma^2 \partial \sigma \circ \circ \\
 & + u' \partial \sigma \circ + 2\Gamma \sigma u' \circ .
 \end{aligned}$$

Remark: Similar to expansions in Feynman diagrams. However, coefficients are not constant but depend on the solution itself.

Higher order expansion

Graphical notations for symbols: \circ for noise, --- for heat kernel, --- for derivative of heat kernel. For example $(\Pi_{z_0}^{\circ})(z) = v(z) - v(z_0)$.
 (! Suitable recentering required !) One has

$$\begin{aligned}
 U = & u \mathbf{1} + \sigma \circ + \sigma \partial \sigma \circ \circ + \Gamma \sigma^2 \circ \circ \circ + u' X \\
 & + 2\Gamma \sigma^2 \partial \sigma \circ \circ \circ + 2\Gamma^2 \sigma^3 \circ \circ \circ \circ + \sigma^3 \partial \Gamma \circ \circ \circ \\
 & + \frac{1}{2} \sigma^2 \partial^2 \sigma \circ \circ \circ + \sigma (\partial \sigma)^2 \circ \circ \circ + \Gamma \sigma^2 \partial \sigma \circ \circ \circ \\
 & + u' \partial \sigma \circ + 2\Gamma \sigma u' \circ \circ .
 \end{aligned}$$

Remark: Similar to expansions in Feynman diagrams. However, coefficients are not constant but depend on the solution itself.

Higher order expansion

Graphical notations for symbols: \circ for noise, --- for heat kernel, --- for derivative of heat kernel. For example $(\Pi_{z_0}^{\circ})(z) = v(z) - v(z_0)$.
 (! Suitable recentering required !) One has

$$\begin{aligned}
 U = & u \mathbf{1} + \sigma \circ + \sigma \partial \sigma \circ \circ + \Gamma \sigma^2 \circ \circ + u' X \\
 & + 2\Gamma \sigma^2 \partial \sigma \circ \circ + 2\Gamma^2 \sigma^3 \circ \circ + \sigma^3 \partial \Gamma \circ \circ \\
 & + \frac{1}{2} \sigma^2 \partial^2 \sigma \circ \circ + \sigma (\partial \sigma)^2 \circ \circ + \Gamma \sigma^2 \partial \sigma \circ \circ \\
 & + u' \partial \sigma \circ + 2\Gamma \sigma u' \circ .
 \end{aligned}$$

Remark: Similar to expansions in Feynman diagrams. However, coefficients are not constant but depend on the solution itself.

Higher order expansion

Graphical notations for symbols: \circ for noise, --- for heat kernel, --- for derivative of heat kernel. For example $(\Pi_{z_0}^{\circ})(z) = v(z) - v(z_0)$.
 (! Suitable recentering required !) One has

$$\begin{aligned}
 U = & u \mathbf{1} + \sigma \text{⦿} + \sigma \partial \sigma \text{⦿}^{\circ} + \Gamma \sigma^2 \text{⦿}^{\circ} + u' X \\
 & + 2\Gamma \sigma^2 \partial \sigma \text{⦿}^{\circ} + 2\Gamma^2 \sigma^3 \text{⦿}^{\circ} + \sigma^3 \partial \Gamma \text{⦿}^{\circ} \\
 & + \frac{1}{2} \sigma^2 \partial^2 \sigma \text{⦿}^{\circ} + \sigma (\partial \sigma)^2 \text{⦿}^{\circ} + \Gamma \sigma^2 \partial \sigma \text{⦿}^{\circ} \\
 & + u' \partial \sigma \text{⦿} + 2\Gamma \sigma u' \text{⦿}.
 \end{aligned}$$

Remark: Similar to expansions in Feynman diagrams. However, coefficients are not constant but depend on the solution itself.

Higher order expansion

Graphical notations for symbols: \circ for noise, $—$ for heat kernel, $—$ for derivative of heat kernel. For example $(\Pi_{z_0}^{\circ})(z) = v(z) - v(z_0)$.
 (! Suitable recentering required !) One has

$$\begin{aligned}
 U = & u \mathbf{1} + \sigma \circ + \sigma \partial \sigma \circ \circ + \Gamma \sigma^2 \circ \circ + u' X \\
 & + 2\Gamma \sigma^2 \partial \sigma \circ \circ + 2\Gamma^2 \sigma^3 \circ \circ + \sigma^3 \partial \Gamma \circ \circ \\
 & + \frac{1}{2} \sigma^2 \partial^2 \sigma \circ \circ + \sigma (\partial \sigma)^2 \circ \circ + \Gamma \sigma^2 \partial \sigma \circ \circ \\
 & + u' \partial \sigma \circ + 2\Gamma \sigma u' \circ .
 \end{aligned}$$

Remark: Similar to expansions in Feynman diagrams. However, coefficients are not constant but depend on the solution itself.

Steps

1. Given function U , when does there exist a **unique** function / distribution u described by $\Pi_z U(z)$ locally “near z ” for every z ? Leads to analogues of Hölder / Sobolev / Besov spaces + “**reconstruction theorem**”.
2. Reformulate required operations for U and obtain analogues to classical embedding theorems / Schauder estimates.
3. Derive **fixed point problem** for U . Solution map continuous function of “model” Π .
4. Show that suitable “renormalised” ways of lifting $\xi^{(\varepsilon)}$ to model $\Pi^{(\varepsilon)}$ converge to limit and characterise limit. Derive “renormalised equation”.

Steps

1. Given function U , when does there exist a **unique** function / distribution u described by $\Pi_z U(z)$ locally “near z ” for every z ? Leads to analogues of Hölder / Sobolev / Besov spaces + “**reconstruction theorem**”.
2. Reformulate required operations for U and obtain analogues to classical embedding theorems / Schauder estimates.
3. Derive **fixed point problem** for U . Solution map continuous function of “model” Π .
4. Show that suitable “renormalised” ways of lifting $\xi^{(\varepsilon)}$ to model $\Pi^{(\varepsilon)}$ converge to limit and characterise limit. Derive “renormalised equation”.

Steps

1. Given function U , when does there exist a **unique** function / distribution u described by $\Pi_z U(z)$ locally “near z ” for every z ? Leads to analogues of Hölder / Sobolev / Besov spaces + “**reconstruction theorem**”.
2. Reformulate required operations for U and obtain analogues to classical embedding theorems / Schauder estimates.
3. Derive **fixed point problem** for U . Solution map continuous function of “model” Π .
4. Show that suitable “renormalised” ways of lifting $\xi^{(\varepsilon)}$ to model $\Pi^{(\varepsilon)}$ converge to limit and characterise limit. Derive “renormalised equation”.

Steps

1. Given function U , when does there exist a **unique** function / distribution u described by $\Pi_z U(z)$ locally “near z ” for every z ? Leads to analogues of Hölder / Sobolev / Besov spaces + “**reconstruction theorem**”.
2. Reformulate required operations for U and obtain analogues to classical embedding theorems / Schauder estimates.
3. Derive **fixed point problem** for U . Solution map continuous function of “model” Π .
4. Show that suitable “renormalised” ways of lifting $\xi^{(\varepsilon)}$ to model $\Pi^{(\varepsilon)}$ converge to limit and characterise limit. Derive “renormalised equation”.

Algebraic structure

Need to understand two things:

- How does $\Pi_{z_0}\tau$ relate to $\Pi_{z_1}\tau$ (**reexpansion**)?
- Which **renormalisation** procedures maintain these relations?

Structure: Two combinatorial Hopf algebras acting on symbols, basis vectors indexed by decorated trees / forests. Reexpansion: generalisation of Connes-Kreimer Hopf algebra / Butcher group. Renormalisation: generalisation of substitution Hopf algebra (analysis of B-series in numerical analysis).

Details depend on problem at hand. Can all be obtained in a functorial way from a single structure where linear maps are replaced by a different kind of morphism.

Algebraic structure

Need to understand two things:

- How does $\Pi_{z_0}\tau$ relate to $\Pi_{z_1}\tau$ (reexpansion)?
- Which renormalisation procedures maintain these relations?

Structure: Two combinatorial Hopf algebras acting on symbols, basis vectors indexed by decorated trees / forests. Reexpansion: generalisation of Connes-Kreimer Hopf algebra / Butcher group. Renormalisation: generalisation of substitution Hopf algebra (analysis of B-series in numerical analysis).

Details depend on problem at hand. Can all be obtained in a functorial way from a single structure where linear maps are replaced by a different kind of morphism.

Algebraic structure

Need to understand two things:

- How does $\Pi_{z_0} \tau$ relate to $\Pi_{z_1} \tau$ (reexpansion)?
- Which renormalisation procedures maintain these relations?

Structure: Two combinatorial Hopf algebras acting on symbols, basis vectors indexed by decorated trees / forests. Reexpansion: generalisation of Connes-Kreimer Hopf algebra / Butcher group. Renormalisation: generalisation of substitution Hopf algebra (analysis of B-series in numerical analysis).

Details depend on problem at hand. Can all be obtained in a functorial way from a single structure where linear maps are replaced by a different kind of morphism.

Conclusions

1. Can build solution theories for **very singular** parabolic SPDEs very similar to “standard” deterministic PDE theory.
2. Analogues of many classical function spaces and results in this context. (Sobolev, Besov, Hölder, embeddings, etc) **Rule of thumb**: reconstruction operator only requires positive “regularity index”.
3. Some **missing pieces**: general proof of convergence of “models”, general “central limit theorem” anywhere near optimality.

Conclusions

1. Can build solution theories for **very singular** parabolic SPDEs very similar to “standard” deterministic PDE theory.
2. Analogues of many classical function spaces and results in this context. (Sobolev, Besov, Hölder, embeddings, etc) **Rule of thumb**: reconstruction operator only requires positive “regularity index”.
3. Some **missing pieces**: general proof of convergence of “models”, general “central limit theorem” anywhere near optimality.

Conclusions

1. Can build solution theories for **very singular** parabolic SPDEs very similar to “standard” deterministic PDE theory.
2. Analogues of many classical function spaces and results in this context. (Sobolev, Besov, Hölder, embeddings, etc) **Rule of thumb**: reconstruction operator only requires positive “regularity index”.
3. Some **missing pieces**: general proof of convergence of “models”, general “central limit theorem” anywhere near optimality.